

FREE OSCILLATION OF WATER IN A LAKE AND THE I -FUNCTION ASSOCIATED WITH THE M-SERIES

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Abstract- The object of this paper is to establish an infinite integral involving the product of M -series and the I -function. We have also obtained a solution of a problem of free oscillation of water in a circular lake with the help of the integral.

1. Introduction

The I -function which was introduced by Saxena [6] is an extension of Fox's H -function. On specializing the parameters this I -function reduces to Meijer's G -function, MacRobert's E -function, Lauricella function, Appell functions, Whittakar function etc. The importance of I -function lies in the fact that almost all the elementary and special functions in the literature follow as its special cases, therefore, the results established in this paper are of a general character and hence encompass several cases of interest.

We will represent and define here the I -function in the following manner:

$$I_{p_i, q_i; r}^{m, n} \left[z \left| \begin{matrix} (a_j, \alpha_j)_{1, n}; (a_{ji}, \alpha_{ji})_{n+1, p_i} \\ (b_j, \beta_j)_{1, m}; (b_{ji}, \beta_{ji})_{m+1, q_i} \end{matrix} \right. \right] = \frac{1}{2\pi i} \int_{\mathcal{L}} \phi(s) x^s ds, \quad (1.1)$$

where

$$\phi(s) = \frac{\prod_{j=1}^m \Gamma(b_j - \beta_j s) \prod_{j=1}^n \Gamma(1 - a_j + \alpha_j s)}{\sum_{i=1}^r \left\{ \prod_{j=m+1}^{q_i} \Gamma(1 - b_{ji} + \beta_{ji} s) \prod_{j=n+1}^{p_i} \Gamma(a_{ji} - \alpha_{ji} s) \right\}}, \quad (1.2)$$

$m, n, p_i (i = 1, \dots, r)$; and $q_i (i = 1, \dots, r)$ are integers satisfying $0 \leq n \leq p_i, 1 \leq m \leq q_i (i = 1, \dots, r)$; r is finite $\alpha_j, \beta_j, \alpha_{ji}, \beta_{ji}$ are real and positive numbers; a_j, b_j, a_{ji}, b_{ji} are complex numbers such that

$$\alpha_j(\beta_h + \nu) \neq \beta_h(\alpha_j - \lambda - 1)$$

for $\lambda, \nu = 0, 1, 2, \dots$; $h = 1, 2, \dots, m$; $i = 1, 2, \dots, r$. \mathcal{L} is a contour running from $\sigma - i\infty$ to $\sigma + i\infty$ (σ is real), in the complex ξ -plane such that the points $s = \frac{(\alpha_j - \lambda - 1)}{\alpha_j}, j =$

$1, 2, \dots, n$; $\lambda = 0, 1, 2, \dots$ and $\xi = \frac{(b_j + \lambda)}{\beta_j}$, $j = 1, 2, \dots, m$; $\lambda = 0, 1, 2, \dots$ lie to the left hand and right hand sides of \mathcal{L} respectively. For the convergence conditions, existence of various contours L and other properties, we can refer Saxena [6, p. 35].

The well known M -series which is a particular case of H -function introduced by Inayat Hussain [3] and is defined by means of the following series expansion.

$${}_u M_v^\tau \left[\begin{matrix} (A_1, \dots, A_u) \\ (B_1, \dots, B_v) \end{matrix}; x \right] = \sum_{k=0}^{\infty} \frac{(A_1)_k \dots (A_u)_k x^k}{(B_1)_k \dots (B_v)_k \Gamma(\tau k + 1)}, \quad (1.3)$$

provided that $\tau \in \mathbb{C}$, $\text{Re}(\tau) > 0$, $(A_j)_k (B_j)_k$ are pochhammer symbols.

2. Formula Required

The following formulae will be required in the present work:

From Erdélyi [2, p. 326, eq. (1)]:

$$\int_0^{\infty} x^{s-1} J_\nu(ax) dx = \frac{2^{s-1} \Gamma\left(\frac{s}{2} + \frac{\nu}{2}\right)}{a^s \Gamma\left(\frac{\nu}{2} - \frac{s}{2} + 1\right)}, \quad -\text{Re}(\nu) < \text{Re}(s) < \frac{3}{2}. \quad (2.1)$$

From Luke [4, p. 291, (6)]:

$$\int_0^{\infty} x^{-1} J_{n+2m+1}(x) dx = \begin{cases} 0, & \text{if } m \neq n \\ (4n + 2a + 2), & \text{if } m = n \end{cases}. \quad (2.2)$$

3. An Infinite Integral

The integral to be evaluated is

$$\begin{aligned} & \int_0^{\infty} x^{\lambda-1} J_\nu(cx) {}_u M_v^\tau \left[\begin{matrix} (A_1, \dots, A_u) \\ (B_1, \dots, B_v) \end{matrix}; hx^{2\rho} \right] I_{p_i, q_i; r}^{m, n} [zx^{-2\mu}] dx \\ &= \frac{2^{\lambda-1}}{c^\lambda} \sum_{k=0}^{\infty} \left(\frac{2}{c}\right)^{2\rho k} \frac{(A_1)_k \dots (A_u)_k h^k}{(B_1)_k \dots (B_v)_k \Gamma(\tau k + 1)} \\ & \cdot I_{p_i, q_i+2; r}^{m+1, n} \left[z \left(\frac{2}{c}\right)^{-2\mu} \left| \begin{matrix} (a_j, \alpha_j)_{1, p} \\ \left(\frac{\lambda}{2} + \rho k + \frac{\nu}{2}, \mu\right), (b_j, \beta_j)_{1, q}, \left(\frac{\lambda}{2} + \rho k - \frac{\nu}{2}, \mu\right) \end{matrix} \right. \right], \end{aligned} \quad (3.1)$$

provided that λ and μ are positive integers, and $\delta > 0, \Omega > 0$, where

$$\Omega = \sum_{j=1}^n \alpha_j - \sum_{j=n+1}^{p_i} \alpha_{ji} + \sum_{j=1}^m \beta_j - \sum_{j=m+1}^{q_i} \beta_{ji} > 0, \quad \delta = \sum_{j=1}^{q_i} \beta_{ji} - \sum_{j=1}^{p_i} \alpha_{ji} > 0$$

Proof

To prove (3.1), in the LHS of equation (3.1) replacing M -series as given with equation (1.3) and I -function as Mellin-Barnes type contour integral (1.1), we find that

$$\int_0^{\infty} x^{\lambda-1} J_{\nu}(cx) \sum_{k=0}^{\infty} \frac{(A_1)_k \dots (A_u)_k x^{2\rho k} h^k}{(B_1)_k \dots (B_v)_k \Gamma(\tau k + 1)} \left[\frac{1}{2\pi\omega} \int_{-\omega\infty}^{\omega\infty} \theta(\xi) z^{\xi} x^{-2\mu\xi} d\xi \right] dx. \quad (3.2)$$

Interchanging the order of integration and summation, and evaluating the inner integral with the help of (2.1), we obtain

$$\begin{aligned} & \sum_{k=0}^{\infty} \frac{(A_1)_k \dots (A_u)_k h^k}{(B_1)_k \dots (B_v)_k \Gamma(\tau k + 1)} \left[\frac{1}{2\pi\omega} \int_{-\omega\infty}^{\omega\infty} \theta(\xi) z^{\xi} d\xi \right] \\ & \cdot \frac{2^{\lambda+2\rho r-2\mu\xi-1} \Gamma\left(\frac{\lambda}{2} + \rho r - \mu\xi + \frac{\nu}{2}\right)}{c^{\lambda+2\rho r-2\mu\xi-1} \Gamma\left(\frac{\nu}{2} - \frac{\lambda}{2} - \rho r + \mu\xi + 1\right)}. \end{aligned} \quad (3.3)$$

Now applying (1.1), the RHS of the integral (3.1) is obtained.

4. Free Oscillation of Water in a Circular Lake

In this section product of M -series and the I -function has been employed in obtaining a solution of the following partial differential equation of oscillation of water in circular lake [5, p. 45 (I)]:

$$x^2 \frac{\partial^2 y}{\partial x^2} + x \frac{\partial y}{\partial x} + \frac{\partial^2 y}{\partial \theta^2} - K^2 x^2 y = 0, \quad (4.1)$$

where y shows the depth of water surface from its position of equilibrium, $K = \frac{w}{\sqrt{(gd)}}$, the velocity of a free wave of small amplitude in a large expanse of water of depth d and g is the acceleration due to gravity with the following considerations

- (1) In any vibrational mode y varies harmonically with time and y is small enough for its square to be neglected.
- (2) The lake is stationary in space.
- (3) There is no loss of energy.

5. Solution of the problem

The solution of (4.1) to be obtained here is

$$Y(x, \theta, t) = \left(\frac{2}{c}\right)^\lambda \sum_{i=0}^{\infty} \frac{i J_i(cx) \cos(i\theta - \phi_i) \cos(\omega it - \psi_i)}{\cos(\phi_i) \cos(\psi_i)} \sum_{k=0}^{\infty} \left(\frac{2}{a}\right)^{2\rho k} \frac{(A_1)_k \dots (A_u)_k}{(B_1)_k \dots (B_v)_k} \\ \cdot \frac{h^k}{\Gamma(\tau k + 1)} I_{p_i, q_i+2; r}^{m+1, n} \left[z \left(\frac{2}{c}\right)^{-2\mu} \left| \begin{array}{c} (a_j, \alpha_j)_{1, p} \\ \left(\frac{\lambda}{2} + \rho k + \frac{\nu}{2}, \mu\right), (b_j, \beta_j)_{1, q}, \left(\frac{\lambda}{2} + \rho k - \frac{\nu}{2}, \mu\right) \end{array} \right. \right]. \quad (5.1)$$

The conditions of validity being the same as stated with (3.1).

PROOF

If we take into account all the physical assumptions [5, p. 47Art .2.61], the complete solution of (4.1) [5, p.47 (2)] is

$$Y(x, \theta, t) = \sum_{i=0}^{\infty} Q_i J_i(cx) \cos(i\theta - \phi_i) \cos(\omega it - \psi_i) \quad (5.2)$$

where $\theta = t = 0$.

Let

$$Y(x, 0, 0) = f(x) = x^\lambda {}_u M_v^\tau \left[\begin{array}{c} (A_1, \dots, A_u); \\ (B_1, \dots, B_v); \end{array} h x^{2\rho} \right] I_{p_i, q_i; r}^{m, n} [z x^{-2\mu}], \quad (5.3)$$

then

$$f(x) = x^\lambda {}_u M_v^\tau \left[\begin{array}{c} (A_1, \dots, A_u); \\ (B_1, \dots, B_v); \end{array} h x^{2\rho} \right] I_{p_i, q_i; r}^{m, n} [z x^{-2\mu}] = \sum_{i=0}^{\infty} Q_i J_i(cx) \cos(\phi_i) \cos(\psi_i). \quad (5.4)$$

Multiplying both sides of (5.4) by $x^{-1} j_\nu(cx)$ integrating with respect to x from 0 to ∞ on the left hand side using (3.1) and on the right hand side using orthogonal property of Bessel function (2.2), we get

$$Q_\nu = \left(\frac{2}{c}\right)^\lambda \frac{\nu}{\cos(\phi_i) \cos(\psi_i)} \sum_{k=0}^{\infty} \left(\frac{2}{a}\right)^{2\rho k} \frac{(A_1)_k \dots (A_u)_k}{(B_1)_k \dots (B_v)_k} \frac{h^k}{\Gamma(\tau k + 1)} \\ \cdot I_{p_i, q_i+2; r}^{m+1, n} \left[z \left(\frac{2}{c}\right)^{-2\mu} \left| \begin{array}{c} (a_j, \alpha_j)_{1, p} \\ \left(\frac{\lambda}{2} + \rho k + \frac{\nu}{2}, \mu\right), (b_j, \beta_j)_{1, q}, \left(\frac{\lambda}{2} + \rho k - \frac{\nu}{2}, \mu\right) \end{array} \right. \right]. \quad (5.5)$$

From (5.2) and (5.5), we obtain the desired result (5.1).

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